

RESEARCH ARTICLE

Novel pathway for mutagenic tautomerization of classical A·T DNA base pairs *via* sequential proton transfer through quasi-orthogonal transition states: A QM/QTAIM investigation

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Citation: Brovarets' OO, Tsiupa KS, Hovorun DM (2018) Novel pathway for mutagenic tautomerization of classical A·T DNA base pairs *via* sequential proton transfer through quasi-orthogonal transition states: A QM/QTAIM investigation. PLoS ONE 13(6): e0199044. <https://doi.org/10.1371/journal.pone.0199044>

Editor: Dennis Salahub, University of Calgary, CANADA

Received: February 11, 2018

Accepted: May 30, 2018

Published: June 27, 2018

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Data Availability Statement: All relevant data are presented within the paper.

Funding: Brovarets' expresses sincere gratitude to organizing committee headed by Prof. Karl Kuchler (Medical University Vienna, Austria) for the financial support (ABC fellow) of her participation in the 7th FEBS Special Meeting "ATP-Binding Cassette (ABC) Proteins: From Multidrug Resistance to Genetic Disease" (March 6–12, 2018, Innsbruck, Austria) and to Chemistry

Abstract

In this paper we have theoretically predicted a novel pathway for the mutagenic tautomerization of the classical A·T DNA base pairs in the free state, the Watson-Crick A·T(WC), reverse Watson-Crick A·T(rWC), Hoogsteen A·T(H) and reverse Hoogsteen A·T(rH) pairs, *via* sequential proton transfer accompanied by a significant change in the mutual orientation of the bases. Quantum-mechanical (QM) calculations were performed at the MP2/aug-cc-pVDZ//B3LYP/6-311++G(d,p) level in vacuum phase, along with Bader's quantum theory of Atoms in Molecules (QTAIM). These processes involve transition states (TSs) with quasi-orthogonal structures (symmetry C_1), which are highly polar, tight ion pairs (A^+ , $N6H_2$ -deprotonated)·(T^+ , $O4/O2$ -protonated). Gibbs free energies of activation for the $A\cdot T(WC) / A\cdot T(rWC) \leftrightarrow A^*\cdot T(rw_{WC}) / A^*\cdot T(w_{WC})$ tautomeric transitions (~ 43.5 kcal·mol⁻¹) are lower than for the $A\cdot T(H) / A\cdot T(rH) \leftrightarrow A^*_{N7}\cdot T(rw_H) / A^*_{N7}\cdot T(w_H)$ tautomerisations (~ 53.0 kcal·mol⁻¹) (rare tautomers are marked by an asterisk; w—wobble configured tautomerisation products). The (T)N3⁺H· · · N1⁻(A), (T)O4⁺H· · · N1⁻(A) / (T)N3⁺H· · · N1⁻(A) and (T)O2⁺H· · · N1⁻(A) H-bonds are found in the transition states $TS^{A^+\cdot T^+}_{A\cdot T(WC) \leftrightarrow A^*\cdot T(rw_{WC})} / TS^{A^+\cdot T^+}_{A\cdot T(rWC) \leftrightarrow A^*\cdot T(w_{WC})}$. However, in the transition state $TS^{A^+\cdot T^+}_{A\cdot T(H) \leftrightarrow A^*_{N7}\cdot T(rw_H)} / TS^{A^+\cdot T^+}_{A\cdot T(rH) \leftrightarrow A^*_{N7}\cdot T(w_H)}$, the (T)N3⁺H· · · N7⁻(A), (T)O4⁺H· · · N7⁻(A) / (T)N3⁺H· · · N7⁻(A) and (T)O2⁺H· · · N7⁻(A) H-bonds are supplemented by the attractive (T)O4⁺/O2⁺· · · N6⁻(A) van der Waals contacts. It was demonstrated that the products of the tautomerization of the classical A·T DNA base pairs— $A^*\cdot T(rw_{WC})$, $A^*_{N7}\cdot T(rw_H)$ and $A^*_{N7}\cdot T(w_H)$ (symmetry C_s)—further transform *via* double proton transfer into the energetically favorable wobble $A\cdot T^*(rw_{WC})$, $A\cdot T^*(rw_H)$ and $A\cdot T^*_{O2}(w_H)$ base mispairs (symmetry C_s).

Biological Interface Division of the Royal Society of Chemistry (RSC, UK) for the RSC Travel Grant for the participation at the “3rd Green and Sustainable Chemistry Conference” (May 13-16, 2018, Hotel Intercontinental, Berlin, Germany). The funders had no role in study design, data collection and analysis, decision to publish, or preparation of the manuscript.

Competing interests: The authors have declared that no competing interests exist.

Introduction

Investigation of microstructural mechanisms for mutagenic tautomerization of the Watson-Crick DNA base pairs occupies an important place in molecular biophysics and molecular biology, enabling an understanding of the nature of genome instability [1–5]. This follows from the ‘rare tautomer hypothesis’ proposed by Watson and Crick [1] shortly after they established the spatial architecture of DNA [2]. However, achievements in this area remain rather modest despite its long history [6], encouraging further research in this direction.

Löwdin [3, 4] first proposed that the electronic structure of the Watson-Crick (WC) DNA base pairs A·T(WC) and G·C(WC) permits their transition into the high-energy tautomerized states A*·T*(L) and G*·C*(L), now called Löwdin (L) base pairs. Here and henceforth, rare (in particular mutagenic) tautomers are marked with an asterisk and differ from each other by the location of a particular proton: in the A* rare tautomer proton bonds to N1 nitrogen atom, A*_{N7} –to N7 nitrogen atom; T* –to O4 oxygen atom and T*_{O2} –to O2 oxygen atom. Löwdin proposed that the A·T(WC) ↔ A*·T*(L) and G·C(WC) ↔ G*·C*(L) transitions occur by double proton transfer (DPT) along neighboring intermolecular hydrogen (H) bonds *via* proton tunneling. These ideas have been prominent in the field of quantum biology and attracted much theoretical study of the mechanisms of spontaneous transitions and transversions arising during DNA replication [7–14].

Recently, it has become clear that Löwdin’s mechanism does not provide the generation of sufficiently long-lived mutagenic tautomers of the DNA bases, which escape from the replicative DNA-polymerase transforming into their canonical tautomeric forms. The root cause of this observation is the absence of the reverse barrier of tautomerization $\Delta\Delta G$ in the A·T(WC) DNA base pair and its small value in comparison with kT (0.62 kcal·mol⁻¹ under normal conditions) for the G·C(WC) DNA base pair [8, 9, 15–18].

In previous papers [19–27] we proposed an alternative mechanism for mutagenic tautomerization of the A·T(WC) and G·C(WC) base pairs into the corresponding wobble base mispairs and *vice versa*, which mechanism obviates the above difficulties. The chief difference of our mechanism from the Löwdin mechanism is that, in the process of mutagenic tautomerization through sequential proton transfer, the DNA bases shift laterally relative each other into the DNA minor or major grooves, leading to the wobble configuration which contains the mutagenic tautomers [19]. Moreover, it turned out that a similar mechanism works also for the mutagenic tautomerization of purine-purine [21], pyrimidine-pyrimidine [22, 23] and purine-pyrimidine [24–27] DNA base mispairs, which are active players in the field of spontaneous point mutagenesis.

This allows us to assume that it is the intrapair tautomeric transition of the wobble pairs from the main tautomeric form into the rare one with a WC configuration or close to it, and *vice versa*, which is the key to understanding the microstructural mechanisms for spontaneous transitions and transversions during DNA biosynthesis [19–27]. Theoretical analyses of such mechanisms have been experimentally confirmed in part for the A·C(w) and G·T(w) purine-pyrimidine pairs [28–31].

This paper uses QM/QTAIM methods to explore new pathways for mutagenic tautomerization of the classical Watson-Crick A·T(WC), reverse Watson-Crick A·T(rWC), Hoogsteen A·T(H) and reverse Hoogsteen A·T(rH) base pairs with a remarkable biological meaning (for more details, see Refs. [32–49]). These are controlled by transition states with a quasi-orthogonal structure (symmetry C₁) which are highly polar tight ion pairs (A⁻, N6H₂-deprotonated)·(T⁺, O4/O2-deprotonated).

Computational methods

The geometries of all the investigated DNA base pairs and transition states (TSs) were optimized using the Gaussian'09 package [50]. The B3LYP/6-311++G(d,p) level of theory [51–55] was used. This level of theory has successfully proved itself for calculations of similar systems [56–63]. The study included harmonic frequency calculations (using a scaling factor of 0.9668 [64–66]) and intrinsic reaction coordinate (IRC) analysis in the forward and reverse directions from each TS using a Hessian-based predictor-corrector integration algorithm [67] at the B3LYP/6-311++G(d,p) level of theory successfully applied in the previous studies [16, 17, 68, 69]. Local minima and TSs (localized by the synchronous transit-guided quasi-Newton method [70]) were confirmed as such by the absence or presence, respectively, of one imaginary frequency. Standard TS theory was applied to estimate the activation barriers for the tautomerisation reactions [71]. Single point electronic energy calculations were performed for the B3LYP geometries at the MP2/aug-cc-pVDZ level of theory [72, 73]. MP2 has been successfully applied to gain chemical information about similar proton transfer reactions in DNA systems [74–79]. The choice of the MP2 level of theory is caused by the insignificant errors in comparison with CCSD(T) method, that was convincingly shown in the benchmark works of Hobza and Šponer [80, 81].

We have performed investigations for the isolated H-bonded pairs of nucleotide bases, that adequately reflects the processes occurring in real duplex environment [14, 30, 31]. At this we relied on experience received in the previous works [82–85] devoted to related topics and systems, where the negligibly small impact of the stacking and sugar-phosphate backbone on the tautomerisation processes has been shown.

The Gibbs free energy G for all structures was obtained in the following way:

$$G = E_{el} + E_{corr} \quad (1)$$

where E_{el} = electronic energy, while E_{corr} = thermal correction to Gibbs free energy.

Electronic interaction energies ΔE_{int} were calculated at the MP2/6-311++G(2df,pd) level of theory as the difference between the total energy of the base pair and energies of the monomers and corrected for the basis set superposition error (BSSE) [86,87] through the counterpoise procedure [88,89] without consideration of the deformation energies of the monomers due to their relatively small values [90].

Bader's quantum theory of Atoms in Molecules (QTAIM) [91–96] was applied to analyse the electron density distribution, using the AIMAll package [97] for the wave functions obtained at the B3LYP/6-311++G(d,p) level of theory. Presence of a bond critical point (BCP), namely, the so-called (3,-1) BCP, and a bond path between non-covalently connected atoms, as well as a positive value of the Laplacian at this BCP ($\Delta\rho > 0$), were considered as criteria for formation of an H-bond or attractive van der Waals contact [98–100].

The energies of the attractive van der Waals contacts [101, 102] in TSs for tautomeric transitions of the base pairs were calculated by the empirical Espinosa-Molins-Lecomte (EML) formula [103, 104] based on the electron density distribution at the (3,-1) BCPs of the specific contacts:

$$E = 0.5V(r) \quad (2)$$

where $V(r)$ = value of a local potential energy at the (3,-1) BCP.

Energies of conventional AH...B H-bonds were evaluated by the empirical Iogansen formula [105]:

$$E_{AH...B} = 0.33 \cdot \sqrt{\Delta v - 40}, \quad (3)$$

where $\Delta\nu$ = magnitude of the stretching frequency shift for the AH H-bonded group involved in the AH...B H-bond relative to the unbound group. Partial deuteration was applied in order to avoid the effect of vibrational resonances [106–114].

The atom numbering scheme for the DNA bases is as per convention [108].

Results and discussion

These novel pathways for the mutagenic tautomerization of four biologically important A·T DNA base pairs—Watson-Crick A·T(WC), reverse Watson-Crick A·T(rWC), Hoogsteen A·T(H) and reverse Hoogsteen A·T(rH) [32–49]—are portrayed in Figs 1 and 2, with data entered into Tables 1–3.

Conformers of the A·T base pairs remain plane symmetric structures along the entire IRC of tautomerization. This also holds for base pairs tautomerising *via* proton transfer along intermolecular H-bonds as per currently known mechanisms for mutagenic tautomerization of WC pairs [16, 17, 19, 49].

The A·T(WC) / A·T(rWC) / A·T(H) / A·T(rH) \leftrightarrow A*·T(rw_{WC}) / A*·T(w_{WC}) / A*_{N7}·T(rw_H) / A*_{N7}·T(w_H) tautomerisation reactions occur *via* the initial migration of proton localized at the N6 atom of the N6H₂ amino group, leading to the formation of the A⁺·T⁻ ion pair and significant change of the mutual orientation of the bases within the pair, i.e. mutual transformation of the *cis* / *trans* \leftrightarrow *trans* / *cis*-orientation of the N1H and N9H bonds relative to each other (Fig 1). Our new mechanism is controlled by the TSs having quasi-orthogonal structures (symmetry C₁). Further proton transfers to the N1/N7 nitrogen atom causing the rotation of the base and formation of the terminal wobble base mispair. Each of these tautomeric conversions is followed by the asynchronous DPT along the intermolecular H-bonds in the wobble base mispairs (Fig 2).

In all four cases of the novel A·T(WC) / A·T(rWC) / A·T(H) / A·T(rH) \leftrightarrow A*·T(rw_{WC}) / A*·T(w_{WC}) / A*_{N7}·T(rw_H) / A*_{N7}·T(w_H) tautomerisation reactions, the TSs are highly polar (~ 6.8–12.7 D) tight ion pairs (energy of interaction between bases in the pairs ~117–142 kcal·mol⁻¹) (Table 2). These TSs are (A⁻, N6H₂-deprotonated)·(T⁺, O4/O2-protonated) ion pairs. In the TS^{A⁻·T⁺}_{A·T(WC) \leftrightarrow A*·T(rw_{WC})} / TS^{A⁻·T⁺}_{A·T(rWC) \leftrightarrow A*·T(w_{WC})} transition states of tautomerisation the (T)N3⁺H⁺··N1⁻(A) (13.06 / 13.24) and (T)O4⁺ / O2⁺H⁺··N1⁻(A) (8.85 / 8.82 kcal·mol⁻¹) are observed, while for the TS^{A⁻·T⁺}_{A·T(H) \leftrightarrow A*_{N7}·T(rw_H)} / TS^{A⁻·T⁺}_{A·T(rH) \leftrightarrow A*_{N7}·T(w_H)} transition states, the (T)N3⁺H⁺··N7⁻(A) (8.98 / 8.46) and (T)O4⁺ / O2⁺H⁺··N7⁻(A) (5.18 / 4.38) H-bonds are supplemented by attractive (T)O4⁺/O2⁺··N6⁻(A) (2.58 / 3.71 kcal·mol⁻¹) van der Waals contacts (Fig 1, Table 2). At this, the (T)N3⁺H⁺··N1⁻/N7⁻(A) H-bonds (~ 8.5–13.0 kcal·mol⁻¹) are significantly stronger than other specific contacts with increased ellipticity. The weakest among them are the attractive (T)O4⁺/O2⁺··N6⁻(A) (2.58 / 3.71 kcal·mol⁻¹) van der Waals contacts (Table 2).

All TS_{A*·T(rw_{WC}) \leftrightarrow A·T*(rw_{WC})}, TS_{A*·T(w_{WC}) \leftrightarrow A·T*(w_{WC})}, TS_{A*_{N7}·T(rw_H) \leftrightarrow A·T*(rw_H)} and TS_{A*_{N7}·T(w_H) \leftrightarrow A·T*(w_H)} of the DPT reactions are stabilized by the N6-H-N3 covalent bridge and one-single intermolecular H-bond—N1H··O4 (11.61), N1H··O2 (10.94), N7H··O4 (13.76) and N7H··O4 (12.95 kcal·mol⁻¹), accordingly (Fig 2, Table 2).

The non-canonical CH··O H-bonds [110, 111] have been registered in the initial complexes of the tautomerisation: A·T(WC)–C2H··O2 (0.74), A*·T*(L)–C2H··O2 (0.57), A·T(rWC)–C2H··O4 (0.77), A·T(H)–C8H··O2 (0.83), A·T(rH)–C8H··O4 (0.86 kcal·mol⁻¹), which are characterized by low energies E_{CH...O}, estimated by the Espinose-Molins-Lecomte formula [103, 104], decreased electron-topological parameters (ρ , $\Delta\rho$, 100· ϵ) and angles (\angle AH··B), but increased intermolecular distances ($d_{C...O}$ and $d_{H...O}$) in comparison with the canonical H-bonds (for more details see Table 2).

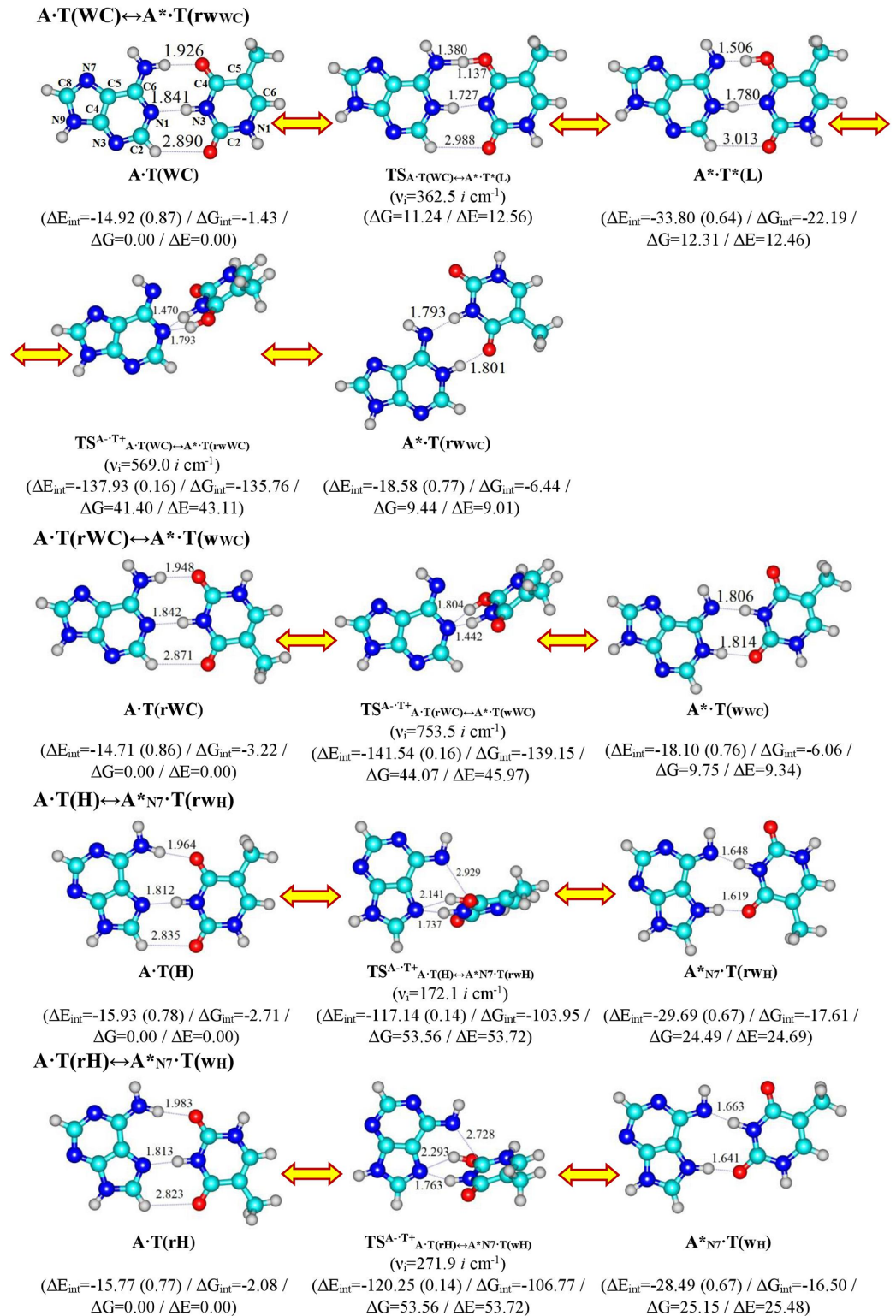


Fig 1. Geometrical structures of the stationary points on the discovered pathways of the tautomerizations via the sequential proton transfer in the four biologically important A·T DNA base pairs through the TSs with quasi-orthogonal oriented bases. Electronic ΔE_{int} (contribution of the total energy of the intermolecular specific contacts) and Gibbs free ΔG_{int} energies of the interaction (MP2/6-311++G(2df,pd)//B3LYP/6-311++G(d,p) level of theory, in kcal·mol⁻¹), relative Gibbs free energies ΔG and electronic energies ΔE (in kcal·mol⁻¹), imaginary frequencies ν_i at the TSs of the

conformational transitions (MP2/aug-cc-pVDZ//B3LYP/6-311++G(d,p) level of theory in the continuum with $\epsilon = 1$ at $T = 298.15$ K) are presented below complexes in brackets. Dotted lines indicate AH...B H-bonds and attractive A...B van der Waals contacts—their lengths H...B and A...B are presented in angstroms (for their more detailed physico-chemical characteristics see Table 2); carbon atoms are in light-blue, nitrogen—in dark-blue, hydrogen—in grey and oxygen—in red.

<https://doi.org/10.1371/journal.pone.0199044.g001>

In general, the values of the electron density ρ at the (3,-1) BCPs of the intermolecular H-bonds range from 0.013 a.u. up to the 0.107 a.u.; the values of the Laplacian of the electron density $\Delta\rho$ at the (3,-1) BCPs are positive for all intrapair H-bonds and lie within a wide range from 0.005 a.u. up to the 0.152 a.u., demonstrating that H-bonds are attractive closed-shell interactions; the value of the ellipticity ϵ varies in the range $0.79\text{--}8.6 \cdot 10^{-3}$ (Table 2).

The classical geometrical criteria are satisfied for all canonical H-bonds in the investigated base mispairs and TSs of their interconversions: $d_{A...B}$ (2.574–3.103 Å), $d_{H...B}$ (1.442–2.293 Å) and $\angle AH...B$ (139.2–179.3°) (Table 2).

Interestingly, the energy of the intermolecular specific contacts (H-bonds and attractive van der Waals contacts) constitute only a minor part of the electronic energy of monomeric interactions for all these H-bonded structures (~14–0.87%) (see Figs 1 and 2). This agrees with previous results for other H-bonded base pairs [112].

All tautomeric transitions in this work are dipole-active, being accompanied by significant changes in dipole moment of the tautomerizing structures along the IRC (0.38–12.65 D), achieving maximum values for each tautomeric transition at its TS (7.38, 6.83, 12.65 and 10.636 D, accordingly) (Table 2). The Gibbs free energy of activation for the A·T(WC)/A·T(rWC) \leftrightarrow A*·T(rw_{WC})/A*·T(w_{WC}) tautomerisations (~ 43.5 kcal·mol⁻¹) is noticeably lower than for the A·T(H)/A·T(rH) \leftrightarrow A·T*(rw_H)/A·T*(o₂(w_H)) tautomerisations (~ 53.0 kcal·mol⁻¹) (Figs 1 and 2).

Note that only one case of mutagenic tautomerization, the A·T(WC) \leftrightarrow A*·T(rw_{WC}) reaction, occurs by participation of the dynamically unstable intermediate A*·T*(L) (a Löwdin's base pair [3, 4]). The other three A·T DNA base pairs—A·T(rWC), A·T(H) and A·T(rH)—do not tautomerise *via* the Löwdin's mechanism. For these three pairs, the local minima corresponding to the tautomerized A*·T*(o₂), A*_{N7}·T* and A*_{N7}·T*(o₂) base pairs are absent on the energy hypersurface. This observation is independent of the level of QM theory used.

It should be noted that three out of four tautomerization processes of the A·T base pairs do not complete with formation of the A*·T(rw_{WC}), A*_{N7}·T(rw_H) and A*_{N7}·T(w_H) mispairs (Fig 2 and Table 1). These plane-symmetric wobble pairs (symmetry C_s) tautomerise further *via* the DPT mechanism along neighboring intermolecular H-bonds into the energetically-favorable plane-symmetric A·T*(rw_{WC}), A·T*(rw_H) and A·T*(o₂(w_H)) DNA base mispairs, respectively (Fig 2, Tables 1 and 2). These processes occur *via* a concerted asynchronous mechanism with proton transfer along the intermolecular (T)N3H...N6(A) H-bonds, which, in fact, is a rate-limiting stage. It is noteworthy that the A*_{N7}·T(rw_H) \rightarrow A·T*(rw_H) and A*_{N7}·T(w_H) \rightarrow A·T*(o₂(w_H)) tautomerisations are barrier-less ($\Delta\Delta G_{TS} = -1.98$ and -1.97 kcal·mol⁻¹) (Table 1), while the activation barriers for the A*·T(rw_{WC}) \leftrightarrow A·T*(rw_{WC}) (1.39) and A*·T(w_{WC}) \leftrightarrow A·T*(o₂(w_{WC})) (1.77) are significantly lower than for the novel tautomerisation reactions (41.40–53.56 kcal·mol⁻¹), but are comparable with the values for the other DPT reactions [113]: from 2.42 for A*·G*_{syn} \leftrightarrow A·G*_{syn} [100] to 10.29 kcal·mol⁻¹ for A·T \leftrightarrow A*·T* [16] DPT tautomerisations.

It is thus possible to say that the tautomerization processes described here terminate with the mutagenic tautomerization of both T and A DNA bases with further formation of the classical mutagenic tautomers T*, T*(o₂) [16, 19, 27, 64, 102, 106] and A* [16, 19–22, 25, 26, 100, 114], respectively. In this case, the A*_{N7}·T(rw_H) \leftrightarrow A·T*(rw_H) and A*_{N7}·T(w_H) \leftrightarrow A·T*(o₂(w_H)) tautomeric equilibria are completely shifted to the right. For the two other cases, the following

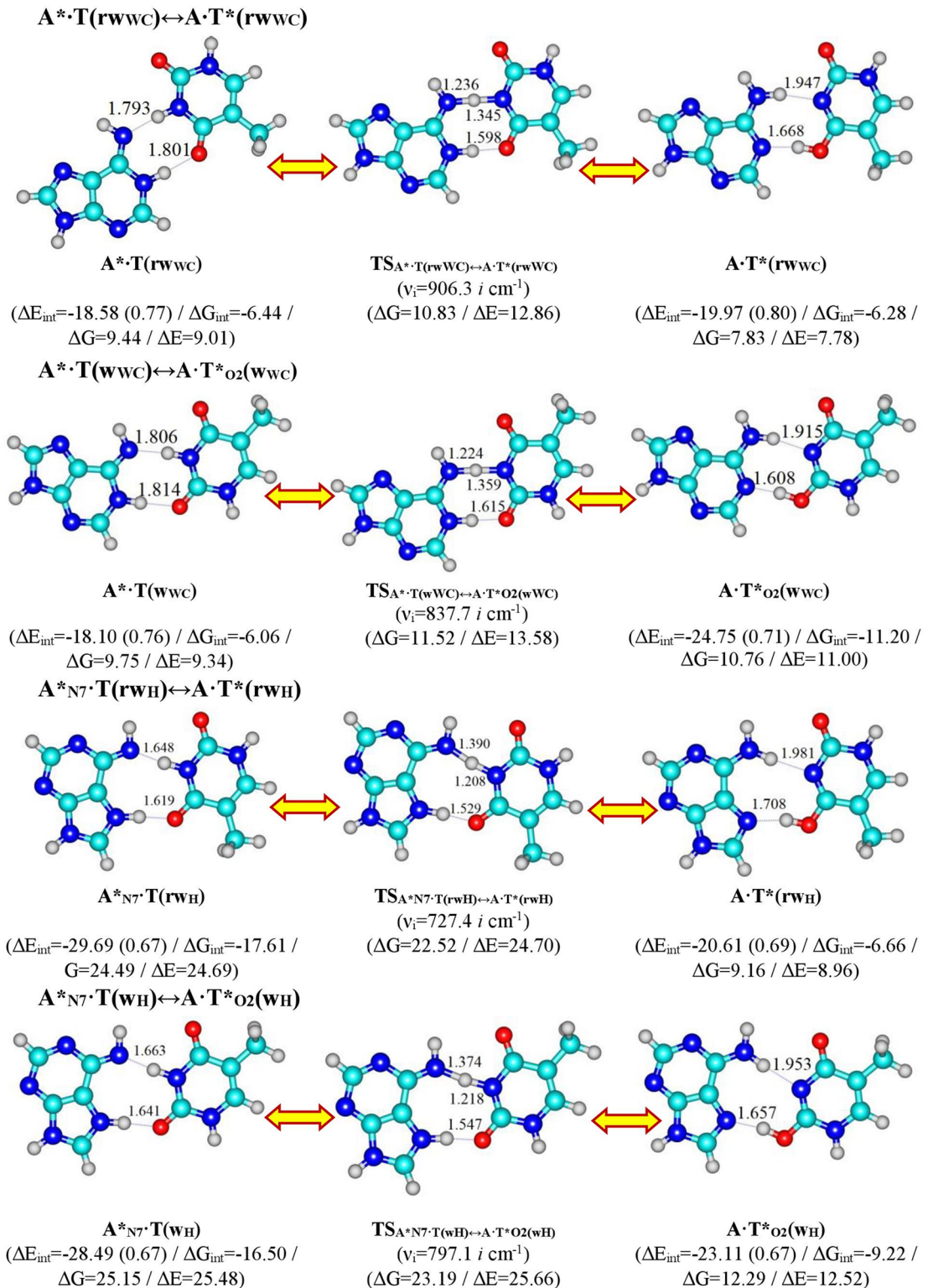


Fig 2. Geometrical structures of the stationary points on the pathways of the tautomerizations via the double proton transfer in the products of the discovered tautomerizations of the classical A·T DNA base pairs. For the detailed designations see Fig 1.

<https://doi.org/10.1371/journal.pone.0199044.g002>

Table 1. Energetic characteristics (in kcal·mol⁻¹) of the discovered mutagenic tautomerizations of the biologically important A·T DNA base pairs via the single and double proton transfers obtained at the MP2/aug-cc-pVDZ//B3LYP/6-311++G(d,p) level of QM theory in the continuum with ε = 1 under normal conditions (see Figs 1 and 2).

| Tautomeric transition | ν_i^a | ΔG^b | ΔE^c | $\Delta\Delta G_{TS}^d$ | $\Delta\Delta E_{TS}^e$ | $\Delta\Delta G^f$ | $\Delta\Delta E^g$ |
|---|-----------|--------------|--------------|-------------------------|-------------------------|--------------------|--------------------|
| A·T(WC) ↔ A*·T*(L) | 362.5 | 12.31 | 12.46 | 11.24 | 12.56 | -1.07 | 0.10 |
| A*·T*(L) ↔ A*·T(rw _{WC}) | 569.0 | 9.44 | 9.01 | 41.40 | 43.11 | 31.95 | 34.10 |
| A·T(rWC) ↔ A*·T(w _{WC}) | 753.5 | 9.75 | 9.34 | 44.07 | 45.97 | 34.32 | 36.63 |
| A·T(H) ↔ A* _{N7} ·T(rw _H) | 172.1 | 24.49 | 24.69 | 52.59 | 52.49 | 28.10 | 27.80 |
| A·T(rH) ↔ A*·T*(rw _H) | 271.9 | 25.15 | 25.48 | 53.56 | 53.72 | 28.41 | 28.24 |
| A*·T(rw _{WC}) ↔ A·T*(rw _{WC}) | 906.3 | -1.61 | -1.24 | 1.39 | 3.85 | 3.00 | 5.08 |
| A*·T(w _{WC}) ↔ A·T* _{O2} (w _{WC}) | 837.7 | 1.01 | 1.66 | 1.77 | 4.24 | 0.76 | 2.58 |
| A* _{N7} ·T(rw _H) ↔ A·T*(rw _H) | 727.4 | -15.34 | -15.73 | -1.98 | 0.01 | 13.36 | 15.73 |
| A* _{N7} ·T(w _H) ↔ A·T* _{O2} (w _H) | 797.1 | -12.87 | -12.96 | -1.97 | 0.18 | 10.90 | 13.14 |

^aImaginary frequency at the TS of the tautomeric transition, cm⁻¹.

^bThe Gibbs free energy of the product relatively the reactant of the tautomeric transition (T = 298.15 K).

^cThe electronic energy of the product relatively the reactant of the tautomeric transition.

^dThe Gibbs free energy barrier for the forward tautomeric transition.

^eThe electronic energy barrier for the forward tautomeric transition.

^fThe Gibbs free energy barrier for the reverse tautomeric transition.

^gThe electronic energy barrier for the reverse tautomeric transition.

<https://doi.org/10.1371/journal.pone.0199044.t001>

proportions are observed: A*·T(rw_{WC}) (6.9%) ↔ A·T*(rw_{WC}) (93.1%) and A*·T(w_{WC}) (83.6%) ↔ A·T*_{O2}(w_{WC}) (16.4%).

Gibbs free energies (in kcal·mol⁻¹) and populations of the investigated base mispairs yield the order: A·T(H) (0.00) < A·T(rH) (0.22/0.63) < A·T(WC) (1.05/0.21) < A·T(rWC) (1.31/0.16) < A·T*(rw_{WC}) (8.83/3.77·10⁻⁷) < A·T*(rw_H) (8.96/1.92·10⁻⁷) < A*·T(rw_{WC}) (10.07/2.47·10⁻⁸) < A*·T(w_{WC}) (10.65/1.13·10⁻⁸) < A·T*_{O2}(w_{WC}) (12.31/2.07·10⁻⁹) < A·T*_{O2}(w_H) (12.74/6.13·10⁻¹⁰) < A*·T*(L) (13.51/1.95·10⁻¹⁰) < A*_{N7}·T(rw_H) (24.69/1.08·10⁻¹⁸) < A*_{N7}·T(w_H) (25.70/2.24·10⁻¹⁹). Notably, populations of the wobble A·T*(rw_{WC}), A·T*(rw_H), A*·T(rw_{WC}), A*·T(w_{WC}), A·T*_{O2}(w_{WC}), A·T*_{O2}(w_H) (12.74/6.13·10⁻¹⁰) and A*·T*(L) tautomerised states, fitting into the range of the frequencies of the spontaneous point mutations observed experimentally (10⁻¹¹–10⁻⁹) [115–117], point on their involvement into the processes of the origin of the spontaneous point mutations.

Notably, the methyl group of the T DNA base does not change its orientation during all these tautomerisation processes without exception. Moreover, the heterocycles of the DNA bases remain planar, despite their ability for out-of-plane bending [118–120].

A relatively small non-planarity of the pyrimidine ring of the protonated T⁺ base occurs only in the TS^{A⁻·T⁺}_{A·T(WC)↔A*·T(rw_{WC})}, TS^{A⁻·T⁺}_{A·T(rWC)↔A*·T(w_{WC})}, TS^{A⁻·T⁺}_{A·T(H)↔A*_{N7}·T(rw_H)} and TS^{A⁻·T⁺}_{A·T(rH)↔A*_{N7}·T(w_H)} transition states. The maximum value of the non-planar dihedral angle reaches 2.5° (C2-N3-C4-C5), 3.1° (N1-C2-N3-C4), 3.7° (C2-N3-C4-C5) and 7.8° (N1-C2-N3-C4), respectively. Another structural feature of the protonated T⁺ base in these TSs is the deviation of the O4⁺H / O2⁺H hydroxyl group from the plane of the pyrimidine ring (the dihedral angles range from 9.3 to 40.1°) (Table 3).

Conclusions and perspectives

Novel pathways for mutagenic tautomerization of four classical A·T DNA base pairs, followed by the significant changes of base orientation within the pair, have been predicted by these QM results. The transition states with quasi-orthogonal structure (symmetry C₁) are highly polar

Table 2. Electron-topological, geometrical and energetic characteristics of the intermolecular specific contacts—H-bonds and attractive van der Waals (vdW) contacts in the investigated DNA base pairs and TSs of their tautomeric transformations obtained at the B3LYP/6-311++G(d,p) level of QM theory ($\epsilon = 1$) (see Figs 1 and 2).

| Complex | AH...B H-bond / A...B vdW contact | ρ^a | $\Delta\rho^b$ | $100\cdot\epsilon^c$ | $d_{A...B}^d$ | $d_{H...B}^e$ | $\angle AH...B^f$ | $E_{AH...B} E_{A...B}^g$ | μ^h |
|---|-------------------------------------|----------|----------------|----------------------|---------------|---------------|-------------------|--------------------------|---------|
| A·T(WC) [16, 19] | N6H...O4 | 0.026 | 0.093 | 4.39 | 2.946 | 1.926 | 173.5 | 4.65 | 1.88 |
| | N3H...N1 | 0.040 | 0.093 | 6.49 | 2.886 | 1.841 | 178.8 | 7.58 | |
| | C2H...O2 | 0.004 | 0.014 | 3.40 | 3.975 | 2.890 | 132.3 | 0.74* | |
| TS _{A·T(WC)→A*·T*(L)} | N1H...N3 | 0.052 | 0.101 | 6.08 | 2.775 | 1.727 | 171.4 | 9.62** | 0.38 |
| | C2H...O2 | 0.004 | 0.012 | 14.48 | 3.722 | 2.988 | 125.3 | 0.61* | |
| A*·T*(L) | O4H...N6 | 0.087 | 0.065 | 4.56 | 2.578 | 1.506 | 174.8 | 13.47 | 0.78 |
| | N1H...N3 | 0.045 | 0.101 | 6.24 | 2.825 | 1.780 | 171.1 | 7.73 | |
| | C2H...O2 | 0.003 | 0.012 | 3.40 | 4.098 | 3.013 | 125.1 | 0.57* | |
| TS ^{A·T+} _{A·T(WC)→A*·T(rwWC)} | N3 ⁺ H...N1 ⁻ | 0.100 | 0.026 | 5.89 | 2.583 | 1.470 | 158.1 | 13.06 | 7.38 |
| | O4 ⁺ H...N1 ⁻ | 0.045 | 0.092 | 10.15 | 2.740 | 1.793 | 154.7 | 8.85 | |
| A*·T(rwWC) | N3H...N6 | 0.044 | 0.095 | 6.22 | 2.844 | 1.793 | 174.7 | 8.53 | 3.23 |
| | N1H...O4 | 0.035 | 0.117 | 3.55 | 2.832 | 1.801 | 177.3 | 5.82 | |
| TS _{A*·T(rwWC)→A*·T*(rwWC)} | N1H...O4 | 0.061 | 0.142 | 3.32 | 2.663 | 1.598 | 179.3 | 11.61** | 3.78 |
| A·T*(rwWC) | N6H...N3 | 0.030 | 0.087 | 7.07 | 2.682 | 1.668 | 170.4 | 5.76 | 2.52 |
| | O4H...N1 | 0.059 | 0.096 | 5.10 | 2.955 | 1.947 | 167.0 | 10.21 | |
| A·T(rWC) [42] | N6H...O2 | 0.024 | 0.088 | 5.26 | 2.962 | 1.949 | 172.9 | 4.38 | 2.40 |
| | N3H...N1 | 0.039 | 0.093 | 6.51 | 2.887 | 1.843 | 177.7 | 7.55 | |
| | C2H...O4 | 0.004 | 0.014 | 3.32 | 3.696 | 2.872 | 132.8 | 0.77* | |
| TS ^{A·T+} _{A·T(rWC)→A*·T*(wWC)} | N3 ⁺ H...N1 ⁻ | 0.107 | 0.005 | 5.60 | 2.574 | 1.442 | 158.2 | 13.24 | 6.83 |
| | O2 ⁺ H...N1 ⁻ | 0.043 | 0.090 | 9.82 | 2.741 | 1.804 | 152.8 | 8.82 | |
| A*·T(wWC) | N3H...N6 | 0.042 | 0.095 | 6.21 | 2.858 | 1.806 | 173.3 | 8.31 | 4.29 |
| | N1H...O2 | 0.034 | 0.115 | 4.40 | 2.845 | 1.814 | 177.0 | 5.49 | |
| TS _{A*·T(wWC)→A*·T*O2(wWC)} | N1H...O2 | 0.058 | 0.141 | 4.10 | 2.676 | 1.615 | 179.2 | 10.94*** | 5.33 |
| A·T*O2(wWC) | N6H...N3 | 0.034 | 0.088 | 1.71 | 2.944 | 1.915 | 167.1 | 6.19 | 3.96 |
| | O2H...N1 | 0.071 | 0.081 | 0.86 | 2.644 | 1.608 | 171.9 | 11.43 | |
| A·T(H) [42] | N6H'...O4 | 0.023 | 0.086 | 3.93 | 2.972 | 1.963 | 170.6 | 4.18 | 6.16 |
| | N3H...N7 | 0.041 | 0.099 | 5.75 | 2.853 | 1.811 | 175.9 | 7.39 | |
| | C8H...O2 | 0.005 | 0.016 | 7.71 | 3.524 | 2.835 | 121.7 | 0.83* | |
| TS ^{A·T+} _{A·T(H)→A*·N7·T(rwH)} | N3 ⁺ H...N7 | 0.050 | 0.097 | 4.51 | 2.754 | 1.737 | 158.7 | 8.98 | 12.65 |
| | O4 ⁺ H...N7 | 0.019 | 0.058 | 9.97 | 3.021 | 2.141 | 147.7 | 5.18 | |
| | O4 ⁺ ...N6 ⁻ | 0.013 | 0.043 | 68.81 | 2.929 | - | - | 2.58* | |
| A* _{N7} ·T(rwH) | N3H...N6 | 0.062 | 0.090 | 5.55 | 2.731 | 1.648 | 174.5 | 11.26 | 9.42 |
| | N7H...O4 | 0.055 | 0.147 | 2.33 | 2.671 | 1.619 | 175.8 | 8.61 | |
| TS _{A*·N7·T(rwH)→A*·T*(rwH)} | N7H...O4 | 0.070 | 0.151 | 2.34 | 2.603 | 1.529 | 175.7 | 13.76** | 8.37 |
| A·T*(rwH) | N6H'...N3 | 0.027 | 0.082 | 7.62 | 3.000 | 1.981 | 175.7 | 5.09 | 7.36 |
| | O4H...N7 | 0.052 | 0.102 | 4.48 | 2.702 | 1.708 | 166.4 | 9.18 | |
| A·T(rH) [42] | N6H'...O2 | 0.022 | 0.082 | 4.95 | 2.994 | 1.986 | 170.9 | 3.90 | 5.67 |
| | N3H...N7 | 0.041 | 0.099 | 5.80 | 2.856 | 1.815 | 176.9 | 7.34 | |
| | C8H...O4 | 0.005 | 0.017 | 7.97 | 3.517 | 2.825 | 121.9 | 0.86* | |
| TS _{A·T(rH)→A*·N7·T(wH)} | N3 ⁺ H...N7 | 0.047 | 0.098 | 3.26 | 2.757 | 1.763 | 154.8 | 8.46 | 10.36 |
| | O2 ⁺ H...N7 | 0.014 | 0.044 | 12.37 | 3.103 | 2.293 | 139.2 | 4.38 | |
| A* _{N7} ·T(wH) | O2 ⁺ ...N6 ⁻ | 0.018 | 0.055 | 78.70 | 2.829 | - | - | 3.71* | 10.35 |
| | N3H...N6 | 0.060 | 0.092 | 5.58 | 2.743 | 1.663 | 175.7 | 10.97 | |
| TS _{A*·N7·T(wH)→A*·T*O2(wH)} | N7H...O4 | 0.051 | 0.145 | 3.17 | 2.689 | 1.641 | 176.3 | 8.09 | 9.46 |
| TS _{A*·N7·T(wH)→A*·T*O2(wH)} | N7H...O4 | 0.067 | 0.152 | 3.15 | 2.615 | 1.547 | 176.3 | 12.95** | 9.46 |

(Continued)

Table 2. (Continued)

| Complex | AH...B H-bond / A...B vdW contact | ρ^a | $\Delta\rho^b$ | $100\cdot\epsilon^c$ | $d_{A...B}^d$ | $d_{H...B}^e$ | $\angle AH...B^f$ | $E_{AH...B} E_{A...B}^g$ | μ^h |
|---|-----------------------------------|----------|----------------|----------------------|---------------|---------------|-------------------|--------------------------|---------|
| A·T ⁺ O ₂ (w _H) | N6H'...N3 | 0.029 | 0.086 | 7.38 | 2.974 | 1.953 | 176.4 | 5.38 | 8.23 |
| | O2H...N7 | 0.059 | 0.100 | 4.48 | 2.664 | 1.657 | 168.0 | 10.16 | |

^aThe electron density at the (3,-1) BCP of the specific contact, a.u.

^bThe Laplacian of the electron density at the (3,-1) BCP of the specific contact, a.u.

^cThe ellipticity at the (3,-1) BCP of the specific contact.

^dThe distance between the A and B atoms of the AH...B / A...B specific contact, Å.

^eThe distance between the H and B atoms of the AH...B H-bond, Å.

^fThe H-bond angle, degree.

^gEnergy of the specific contact, calculated by Iogansen's [105], Espinose-Molins-Lecomte [103, 104] (marked with an asterisk) or Nikolaienko-Bulavin-Hovorun [109] (marked with double asterisk) formulas, kcal·mol⁻¹.

^hThe dipole moment of the complex, D.

<https://doi.org/10.1371/journal.pone.0199044.t002>

Table 3. Selected geometrical parameters, characterizing the non-planarity of the discovered mutagenic tautomerizations of the biologically important A·T DNA base pairs *via* the single and double proton transfers, obtained at the B3LYP/6-311++G(d,p) level of QM theory in the continuum with $\epsilon = 1$.

| TS of tautomerisation | Dihedral angles, degree | |
|---|-------------------------|------------------|
| | (A)N7C5(T)N3C4 | (T)HO4/O2C4/C2N3 |
| TS ^{A-T+} _{A·T(WC)↔A*·T(rwWC)} | 85.0 | -9.3 |
| TS ^{A-T+} _{A·T(rWC)↔A*·T(wWC)} | 60.4 | 11.8 |
| TS ^{A-T+} _{A·T(H)↔A*·N7·T(rwH)} | -99.0 | -19.9 |
| TS ^{A-T+} _{A·T(rH)↔A*·N7·T(wH)} | -119.4 | 40.1 |

<https://doi.org/10.1371/journal.pone.0199044.t003>

tight ion pairs (A⁻, N6H₂-deprotonated)·(T⁺, O4/O2-protonated). The tautomerization products—the A^{*}·T(rw_{WC}), A^{*}_{N7}·T(rw_H) and A^{*}_{N7}·T(w_H) pairs—further transform *via* concerted asynchronous double proton transfer into the energetically favorable wobble A^{*}·T*(rw_{WC}), A^{*}·T*(rw_H) and A^{*}·T*_{O2}(w_H) mispairs (symmetry C_s), respectively. Moreover, it was established in our recent papers, that wobble A^{*}·T(rw_{WC}) base mispair can also be formed from the reverse A·T(rWC) base pair [20], A^{*}·T(wWC) base mispair—from the canonical A·T(rWC) base pair [19] and A^{*}_{N7}·T(w_H) base mispair—from the Hoogsteen A·T(H) base pair [20].

We are currently engaged in elaborating this topic in order to discover biologically important H-bonded nucleobase pairs, for which the mechanism of mutagenic tautomerization plays a key role. Moreover, we suggested that novel mechanism of mutagenic tautomerization presented in this study could lead to the conversion of an anti-parallel DNA helix to a parallel DNA helix. We also consider investigation of these tautomerisation mechanism by the participation of the modified A·T base pairs [121–124] as a task for the future.

Supporting information

S1 Dataset. Cartesian coordinates of the investigated complexes: A·T(WC);

TS_{A·T(WC)↔A*·T*(L)}; A^{*}·T*(L); TSA^{-·T+A·T(WC)↔A*·T(rwWC); A^{*}·T(rw_{WC});}

TS_{A*·T(rwWC)↔A·T*(rwWC)}; A·T*(rw_{WC}); A·T(rWC); TSA^{-·T+A·T(rWC)↔A*·T(wWC);}

A^{*}·T(w_{WC}); TSA^{-·T+A·T(H)↔A*·N7·T(rwH); A·T(H); TSA^{-·T+A·T(H)↔A*·N7·T(rwH);}}

A^{*}_{N7}·T(rw_H); TSA^{-·T+A·T(H)↔A*·N7·T(wH); A·T(H); TSA^{-·T+A·T(H)↔A*·N7·T(wH);}}

A^{*}_{N7}·T(w_H); TSA^{-·T+A·T(H)↔A*·N7·T(wH); A·T*(w_H).}

(DOC)

Acknowledgments

The authors gratefully appreciate technical support and computational facilities of joint computer cluster of SSI “Institute for Single Crystals” of the National Academy of Sciences of Ukraine (NASU) and Institute for Scintillation Materials of the NASU incorporated into Ukrainian National Grid.

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